

# Raman-Studies of the Organic Superconductor $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI

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The resonance Raman spectra of (IAuI)<sup>−</sup>-anions in the organic superconductor  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI as well as in the electrolyte salt n-Bu<sub>4</sub>AuI<sub>2</sub> are investigated. In contrast to  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> in the crystals of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI at low temperature only one Raman line at 120 cm<sup>−1</sup> corresponding to the symmetrical stretching mode of the (IAuI)<sup>−</sup>-anion is observed. No distinct change in the Raman spectra is observed at the superconducting transition.

## Introduction

The superconductivity in organic materials is very sensitive against slight structural changes and disorder. Even very small changes can seriously influence the superconducting transition. In the organic charge transfer-salt  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> [BEDT-TTF = bis-(ethylenedithio)tetrathiofulvalene] below 195 K exists an incommensurate modulation [1] and below 125 K a commensurate superstructure with a pronounced distortion of the I<sub>3</sub><sup>−</sup>-anions [2]. Under ambient pressure the superconducting transition temperature  $T_c$  is 1.3 K [3] but after a particular pressure-temperature cycling procedure superconductivity can be stabilized at ambient pressure even at  $T_c = 8.1$  K [4]. A recent neutron diffraction study has shown that the high- $T_c$  superconducting phase of  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> is completely ordered [5] showing that during the pressure-temperature cycling procedure the distortion of the I<sub>3</sub><sup>−</sup>-anions in  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> is suppressed resulting in an increase of  $T_c$ . The I<sub>3</sub><sup>−</sup>-anions and other polyhalide ions used in organic superconductors are sensitive to distortions and slight changes in the crystal. The Raman-spectra of the polyhalide ions exhibit unusually strong resonance lines. On the other hand, resonance Raman (RR)-

scattering is very sensitive to small distortions and changes of intermolecular forces as demonstrated for the I<sub>3</sub><sup>−</sup>-anions in  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> [6]. In [6] the distortions of the I<sub>3</sub><sup>−</sup>-anions and the structural changes were investigated on the basis of RR-studies of the symmetrical stretching mode of the I<sub>3</sub><sup>−</sup>-anions.

In this letter we report on RR-investigations of the organic superconductor  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI. The superconductivity in  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI-crystals is observed below  $T_c = 4-5$  K [7–10], and superconductivity is a volume property [9]. In contradiction to the geometry of the I<sub>3</sub><sup>−</sup>-anions in  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> below 125 K the (IAuI)<sup>−</sup>-anions in  $\beta$ -(BEDT-DDT)<sub>2</sub>IAuI are linear and symmetric in the whole temperature range below 300 K [7]. The structure of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI is completely ordered. However, during the electrochemical process the (IAuI)<sup>−</sup> can decompose and “impurities” obtained in this way might partially replace (IAuI)<sup>−</sup>-anions in the crystal lattice [10]. As a consequence the onset temperature for superconductivity varies over a broad temperature range. In this paper we report the Raman-spectra of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI-crystals below and above the temperature of the superconducting transition. Since n-Bu<sub>4</sub>AuI<sub>2</sub> is used as electrolyte salt in the electrocrystallization of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI the spectra of this compound were also studied.

The linear, symmetrical ion (I–Au–I)<sup>−</sup> has three fundamental modes of vibration: the symmetric stretching mode  $\nu_1$ , the asymmetric stretching mode  $\nu_2$  and a bending mode  $\nu_3$ . Only the stretching mode  $\nu_1$  is Raman-active, however intermolecular inter-

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actions in the solid state can break the symmetry and all three modes might exhibit Raman-activity.

## Experimental

The single crystals of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI were grown by electrochemical methods in the presence of n-Bu<sub>4</sub>NAuI<sub>2</sub> as a supporting electrolyte similarly as described for the analog I<sub>3</sub><sup>-</sup>-compound [11]. The RR-spectra of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI single crystals were measured with a 90° scattering geometry as a function of temperature using three different lines of an argon-ion laser (4579, 4880 and 5145 Å). The beam was focussed on the surface of the crystal and the laser power was below 10 mW. The crystals of n-Bu<sub>4</sub>NAuI<sub>2</sub> were powdered and compressed with KBr in a ratio 1 : 5.

## Results and Discussion

The RR-spectrum of n-BuNAuI<sub>2</sub> dispersed in a KBr-pellet for  $T = 14$  K and an excitation-wavelength  $\lambda = 4579$  Å is shown in Figure 1. The spectra were not essentially dependent on the temperature and the frequency of the exciting light. The observed peak frequencies at 14 K are 31, 36, 91, 113\*, 124\* and  $154 \text{ cm}^{-1} \pm 2 \text{ cm}^{-1}$  (\* indicates the strong lines); moreover some overtones are seen. Similar Raman-results were obtained for polycrystalline samples and for a sample solved in tetrahydrofuran. The strong doublet (113 and 124 cm<sup>-1</sup>) (Fig. 1) can be assigned to the symmetrical stretching mode of the (IAuI)<sup>-</sup>-anion and the splitting of this mode is due to crystal forces in the n-Bu<sub>4</sub>NAuI<sub>2</sub>-crystal. Our spectral results and proposed assignments disagree with [12], where a band at  $157 \text{ cm}^{-1}$  was the strongest and was assigned to the symmetric stretching mode of the (IAuI)<sup>-</sup>-anion. However, the frequency of the symmetrical stretching mode for (IAuI)<sup>-</sup> and for I<sub>3</sub><sup>-</sup> ( $120 \text{ cm}^{-1}$ ) should be close to each other since the structure of these anions is similar and the Au-atom is located in the middle. Moreover, the correctness of our results and proposed assignment is also supported by the results obtained from the  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI-crystals presented below.

The low energetic part of the RR-spectra of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI single crystals for different tem-

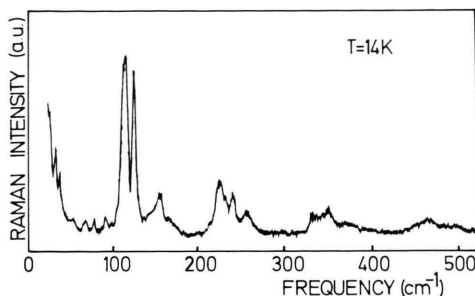


Fig. 1. The Raman spectrum of n-Bu<sub>4</sub>NAuI<sub>2</sub> crystals dispersed in a KBr pellet (weight concentration  $c = 20\%$ , excitation  $\lambda = 4579$  Å).

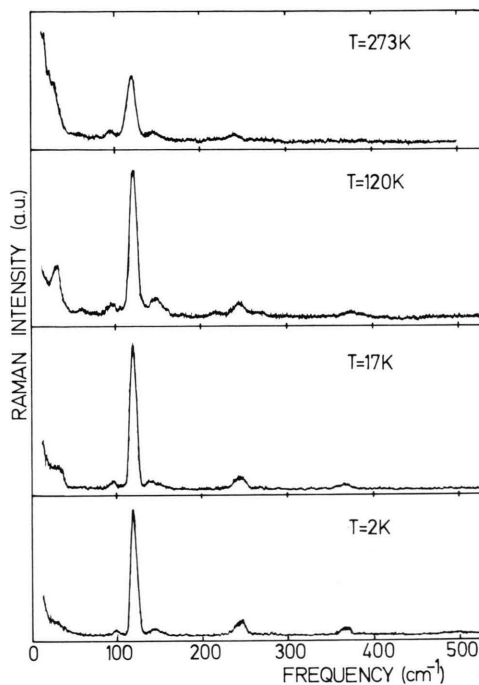


Fig. 2. The temperature dependence of the low energetic part of the resonance Raman spectra of  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI single crystals (excitation  $\lambda = 4579$  Å).

peratures is shown in Fig. 2 (excitation wavelength  $\lambda = 4579$  Å). Similar results were obtained for excitations with light of  $\lambda = 4880$  Å and  $\lambda = 5145$  Å but the intensity of the lines was weaker. The stronger Raman-scattering for an excitation with light of the wavelength of  $\lambda = 4579$  Å is due to the fact that this blue light is closer to the electronic absorption peak of the (IAuI)<sup>-</sup>-anion in the crystals. The RR-bands are observed at about 30, 60, 90, 120\*, 150 and  $175 \text{ cm}^{-1} (\pm 2 \text{ cm}^{-1})$  at a temperature of 120 K (\* indicates the strongest line).

The band at  $120\text{ cm}^{-1}$  can be assigned to the symmetrical stretching mode  $\nu_1$  of the (IAuI)<sup>-</sup>-anions; some overtones of this mode are also seen. The band at  $30\text{ cm}^{-1}$  corresponds to a librational phonon of the (IAuI)<sup>-</sup>-anions. The bands around 90 and  $150\text{ cm}^{-1}$  can be related to the process of creation or annihilation of a symmetric stretching vibration and simultaneous creation of one phonon related to the librational mode ( $30\text{ cm}^{-1}$ ) of (IAuI)<sup>-</sup>. This interpretation is similar to that described in [13] for (BEDT-TTF)<sub>2</sub>X, where X = I<sub>3</sub><sup>-</sup> and IBr<sub>2</sub><sup>-</sup>. The band around  $60\text{ cm}^{-1}$  can be assigned to another lattice mode which also creates the combination band with the symmetric stretching mode  $\nu_1$  ( $175\text{ cm}^{-1}$ ). The band around  $150\text{ cm}^{-1}$  could be related to the asymmetric stretching mode of (IAuI)<sup>-</sup>-anion, but this mode is very weak. The transition to the superconducting state has no distinct influence on the Raman-spectra as is seen from the comparison of the spectra for  $T = 2\text{ K}$  and  $T = 17\text{ K}$  (Figure 2). The band related to the mode  $\nu_1$  is not split even for the best resolution as it was

observed for  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> [6]. This means that no distortion of the IAuI-anions exists, showing that the chains of (IAuI)<sup>-</sup>-anions are regular and every (IAuI)<sup>-</sup>-anion has the same geometry in the metallic as well as in the superconducting state. A disorder due to "impurities" produced during the electrochemical process by decomposition of (IAuI)<sup>-</sup>-ions has negligible influence on the Raman-results since the "impurity"-concentration is small. These "impurities" have an influence on the collective electronic phenomena whereas local phenomena influence the Raman-spectra.

In conclusion, we have performed resonance Raman-studies of the organic superconductor  $\beta$ -(BEDT-TTF)<sub>2</sub>IAuI above and below the superconducting transition temperature. The investigations have shown that the (IAuI)<sup>-</sup>-anions are symmetric and linear in the whole temperature range 1.3–300 K in contrast to the I<sub>3</sub><sup>-</sup>-anions in  $\beta$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> and that no structural change exists below the temperature of the transition to the superconducting state.

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